

Nernst effect

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September 30, 2009

Nernst effect



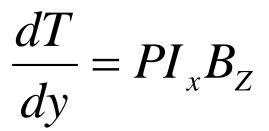
Walther Hermann Nernst

1864-1941

Nernst-Ettingshausen effect (1stNE)

$$E_{y} = N \frac{1}{B} \frac{dT}{dx}$$





Graz 1887 Ludwig Boltzman and coworkers

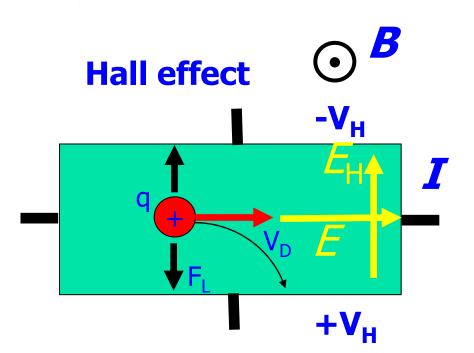
Thesis 1887: Electromotive forces produced by magnetism in heated metal plates

Third law of thermodynamics (Nobel prize 1920) Nernst glower Bernstein-Siemens-Nernst electric piano







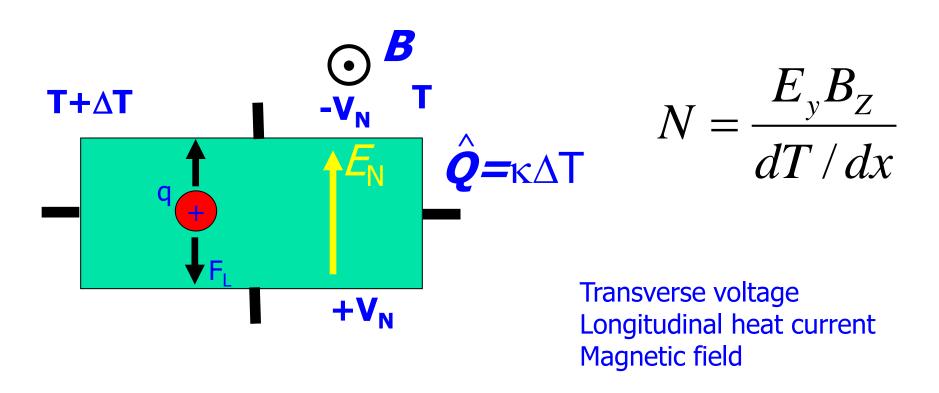


$$V_H = R_H \frac{IB}{d}$$

Transverse voltage Longitudinal current Magnetic field

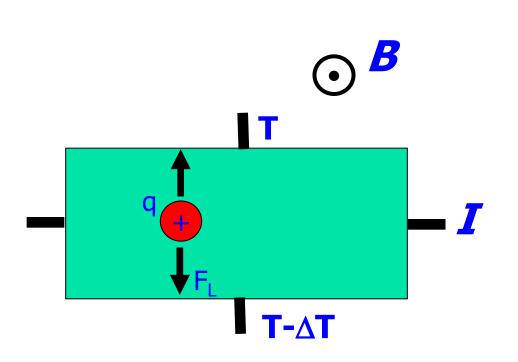


Nernst effect or 1st Nernst-Ettingshausen effect





Ettingshausen effect or 2nd Nernst-Ettingshausen effect



$$\Delta T = P \frac{BI}{d} - \frac{1}{2} \frac{\rho I^2}{d^2 \kappa}$$

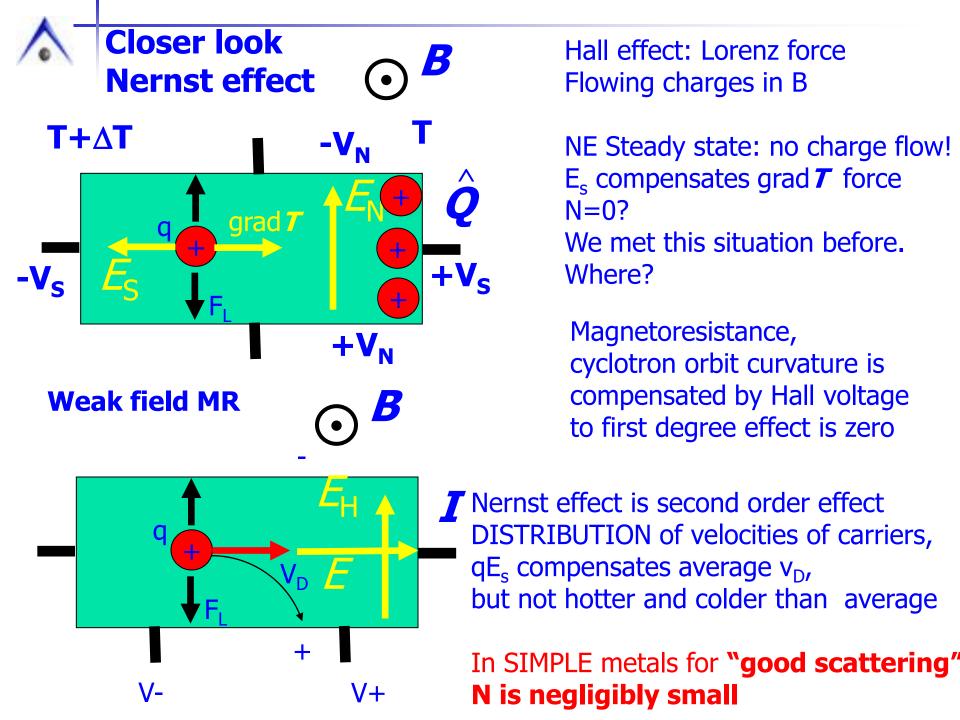
d-sample thickness

 κ - thermal conductivity

ρ- electrical resistivity

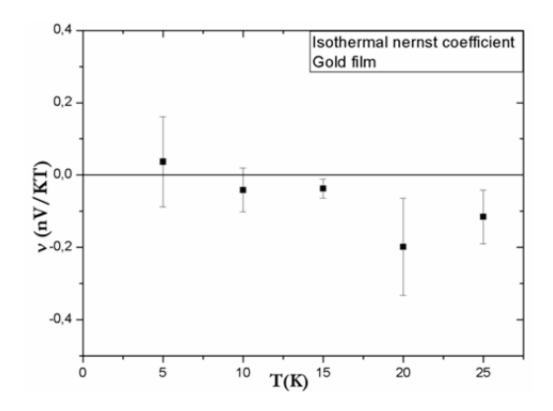
Transverse thermal gradient Longitudinal electrical current Magnetic field

Can be used for thermoelectric cooling





Nernst effect in gold ~0.1 nV/KT





Nernst effect in anisotropic metals

VOLUME 54, NUMBER 9

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"Bad scattering"

Anisotropy notably increases Nernst effect

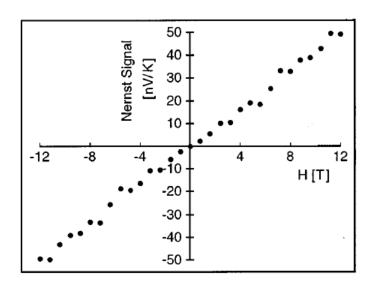


FIG. 2. Magnetic-field dependence of the raw Nernst signal at T=250 K showing the linear variation with the applied field strength.

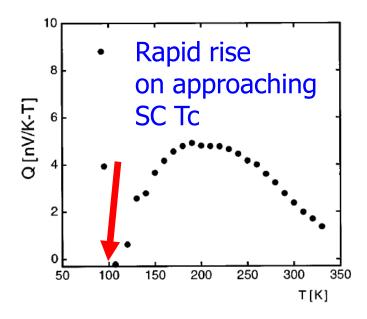


FIG. 3. Temperature dependence of the Nernst coefficient Q of the YBa₂Cu₃O_{7- δ} thin-film sample. Q goes to zero just above T_c , but increases rapidly at lower temperature, reflecting the large contribution from superconducting fluctuations.







Measurement of Nernst effect

Typical numbers

Metals

Good scattering ???

Anisotropic scattering 5 nV/kT

V<10 nV

Semiconductors

 $n\sim10^{16}\text{--}10^{24} \text{ m}^{-3}$ $V\sim1\text{--}100 \text{ }\mu\text{V}$

Very demanding measurements from thermal stability and electrical noise point of view

Ideologically similar to Thermopower measurements

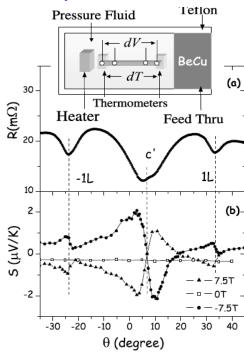


FIG. 1. The experiment setup. (Inset): Four Au wires were attached to the sample along the **a** axis. Two RuO thin film resistors thermometers were placed next to both ends of the sample to measure temperature gradient generated by a miniature heater. The angular dependence of the magnetoresistance (a) and the raw thermoelectric signal (b), $S = \Delta V/\Delta T$, for $(TMTSF)_2PF_6$ as a function of magnetic field direction for $\mathbf{b^*c^*}$ plane rotation (2 K, 10 kbar). $\theta = 0$ is taken as $\mathbf{B} \parallel \mathbf{c^*}$. Note that S is essentially an odd function of applied field. The solid line is a guide to the eye.



Measurements

Nernst signal is defined as odd part of Sxy in field

$$S_{xy} = S_{nonequipotential} + S_{MR} + S_{N}$$

Measurements in positive and negative fields, $Sxy(H)-Sxy(-H)=2S_N$

Fixed temperature +H to -H sweep Time consuming

Strict requirements for T-drift, $Sxy(\delta T) < < Sxy(H)$



Low-frequency method for magnetothermopower and Nernst effect measurements on single crystal samples at low temperatures and high magnetic fields

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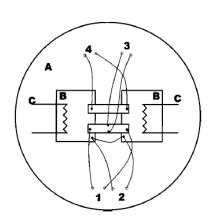


FIG. 1. Diagram of the measurement holder (the outer diameter of the cylindrical copper holder is 10 mm). A: Cu heat sink, B: quartz blocks, and C: heaters. 1: thermopower leads of sample, 2: Chromel–Au(Fe0.07%) thermocouples for ΔT leads, 3: Nernst voltage leads of sample, and 4: thermopower leads of reference YBCO sample.

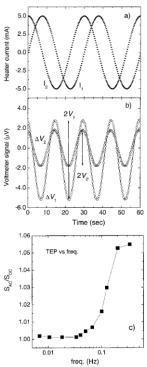


FIG. 2. (a) Heater currents and (b) $\Delta V_1(\Delta V_2)$ as a function of the time. T period of the heating cycle is 30 s and the corresponding periods of oscillation of temperature gradient and thermopower signal are 15 s. (c) S_{w_i}/V_3 is frequency method used to determine the optimum frequency range whe $S_{w_i}/S_{w_i}-1$ for the TEP measurements.

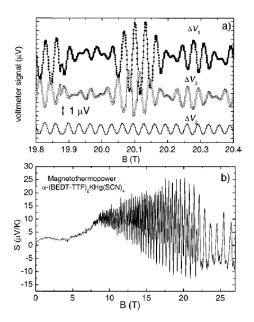


FIG. 4. Magnetothermopower. (a) ΔV_1 , ΔV_2 , and ΔV_3 curves under magnetic field for α -(BEDT-TTF)₂KHg(SCN)₄ at T=0.7 K. (b) Derived magnetothermopower results. Note the narrow range of field in (a), which corresponds to only a few quantum oscillations in (b).



Why bother measuring Nernst effect?

- Additional insight into multiple carrier conductors
- •Anomalous scattering, easy to detect sharp features in $\frac{\partial \sigma(\varepsilon)}{\partial \varepsilon}$
- Exotic scattering in magnetic systems



K. Behnia

•Thermoectricity as a probe of exotic states of correleted electrons is still largely underexplored.



Nernst effect in superconductors

PHYSICAL REVIEW

VOLUME 181, NUMBER 2

10 MAY 1969

Nernst Effect and Flux Flow in Superconductors. I. Niobium*

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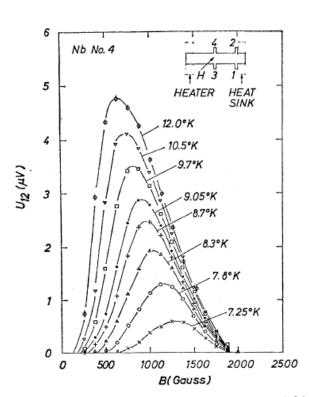


Fig. 2. Transverse voltage U_{12} versus magnetic field for different temperature gradients. The temperature at each curve is the value at the heater. (Specimen 4; temperature at heat sink = 4.2° K.)

N is big in the mixed state of SC Flux quanta respond to grad T by thermal diffusion generate Seebeck and Nernst voltages



Nernst effect in superconductors

Supercond, Sci. Technol, 8 (1995) 189-198. Printed in the UK

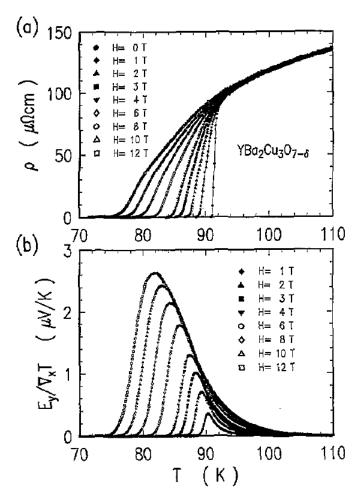


Figure 3. Resistivity ρ (a) and normalized Nernst electric field $E_y/\nabla_x T$ (b) versus temperature for an epitaxial c-axis-oriented YBa₂Cu₃O₇₋₃ film at different magnetic fields ($B \parallel c$).

REVIEW ARTICLE

Superconductors in a temperature gradient

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Same big in high-Tc cuprates But in a broader T-range



Vortex-like excitations and the onset of superconducting phase fluctuation in underdoped $La_{2-x}Sr_xCuO_4$

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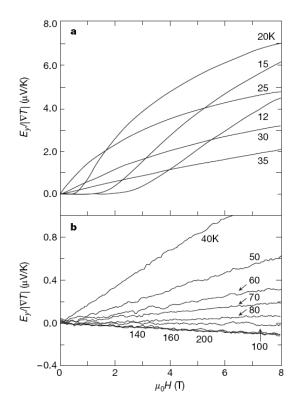


Figure 1 Nernst signals. **a**, The Nernst signal E_y (normalized to unit gradient) versus $\mathbf{H} \parallel \mathbf{c}$ in $\mathrm{La}_{2-x}\mathrm{Sr}_x\mathrm{CuO}_4$ (sample 3, x=0.10) at temperatures $12-35\mathrm{K}$. **b**, The Nernst signal from 40 to 200 K. Above 20 K, the applied gradient is $5\mathrm{K}$ cm⁻¹, while below 20 K, it is half as large. When vortex pinning is large ($T<25\mathrm{K}$), E_y is zero over a range of $H<H_{\mathrm{m}}$. Above 140 K, the curves tend asymptotically to a straight line of negative slope.

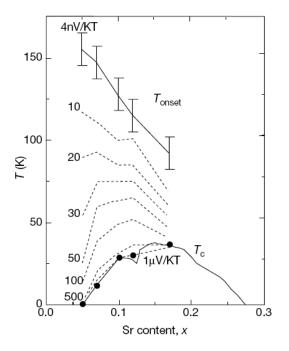


Figure 4 Contour plot of $(\nu-\nu_{\rm n})$ versus x in the phase diagram of LSCO. The contour plot displays how high in ${\cal T}$ the vortex-like excitations extend for each value of x. The upper solid line ${\cal T}_{\rm onset}$ is the contour set by our resolution. The pseudogap ${\cal T}^*$ estimated from heat capacity ¹⁵ is about a factor of two larger than ${\cal T}_{\rm onset}$. Values of ${\cal T}_{\rm c}$ in our samples (circles) match the ${\cal T}_{\rm c}$ line (lower solid line) from Takagi et al. ¹⁴ We note that the ${\cal T}_{\rm c}$ line is roughly similar to the contour line $\nu=1$ μ V/KT.

Claims:

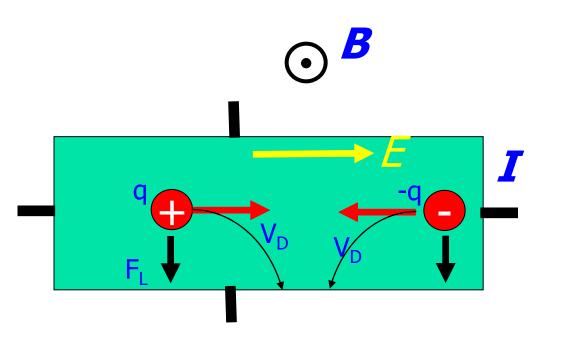
Nernst effect is too big for a metal Superconducting vortices above Tc Preformed pairs scenario

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Can Nernst effect be big otherwise?

To have big MR you need $V_H = 0$



No compensation for cyclotron orbit curvature

Q. What should we have =0 to get big Nernst effect?

A. S = 0
Compensation of different carrier types

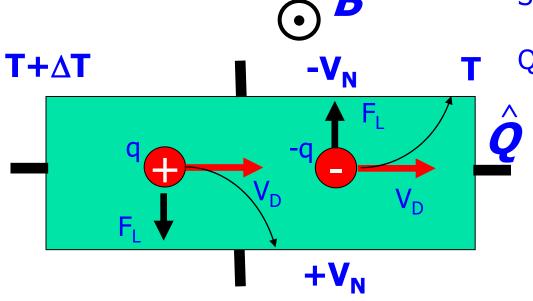
When S=0 there is no restoring force and there is always a current of two carrier types in the same direction



Ambipolar Nernst effect

Contrary to Hall effect
Contributions of +q and -q
Sum up in ambipolar Nernst effect

Q. What is the difference?

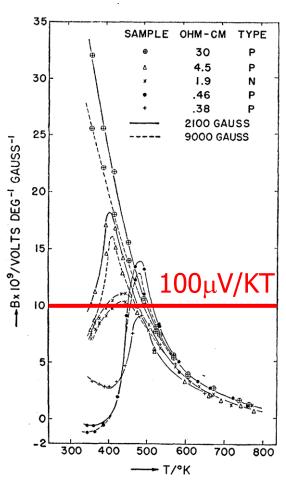


One carrier type: Nernst is second order effect

Two carrier types: Nernst effect is FIRST order effect



Nernst effect Ge: big



Semiconductors
First observation
Krylova TV, Mochan IV, J. Tech.
Phys. (USSR) 25, 2119 (1955)

N is very big in narrow gap semiconductors

Number of unknowns=number of equations To solve transport completely you want to measure Single carrier type: resistivity +Hall

Second carrier type: +Seebeck +Nernst

Fig. 2. Experimental Nernst coefficient, B, in n- and p-type germanium of different resistivities, as a function of temperature between 300 and 750°K, measured at 2100 and 9000 gauss.



Ambipolar Nernst Effect in NbSe₂

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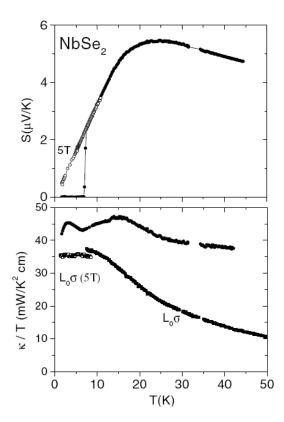


FIG. 2. Upper panel: thermopower (S) of NbSe₂ at H=0 (solid circles) and H=5 T (open circles). Lower panel: thermal conductivity divides by temperature (solid circles) as a function of temperature. Also shown is the charge conductivity (σ) at H=0 (solid squares) and at H=5 T (open squares) multiplied by the constant L_0 (see text).



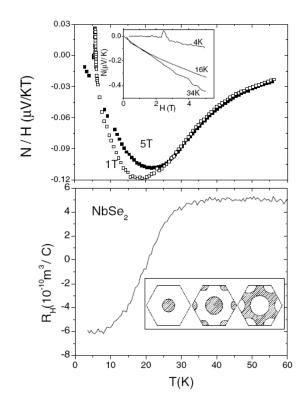


FIG. 3. Upper panel: the Nernst coefficient as a function of temperature at H=0 and H=5 T. The inset compares the field dependence of the Nernst signal at three different temperatures. Lower panel: the temperature dependence of the Hall coefficient measured at H=5 T. Inset: a schematic plot of the three-band Fermi surface in NBSe₂ as observed by angular-resolved photoemission spectroscopy (ARPES) [11].



Nernst effect is big when weak field MR is big: Compensation+ high mobility

Bi metal 200% MR in 2T at room temperature

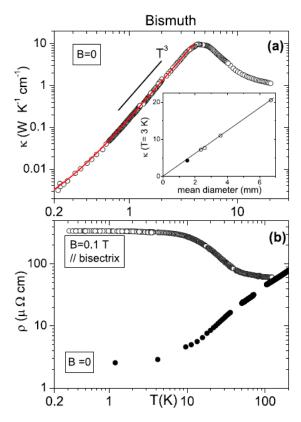


FIG. 1 (color online). (a) Thermal conductivity, κ of the Bi single crystal. Solid line represents a $aT+bT^3$ fit (see text). Inset compares the magnitude of $\kappa(3K)$ of the sample of this study (solid circle) with those reported in Ref. [15] (open circles) as a function of mean diameter. (b) Resistivity of the same sample at zero field and in presence of a field of 0.1 T.

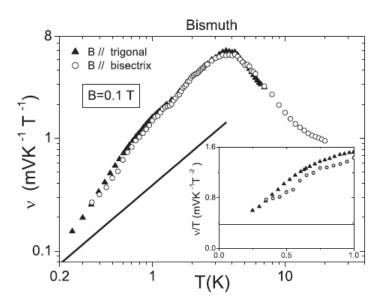


FIG. 2. The temperature dependence of the absolute value of the Nernst coefficient of the bismuth single crystal for two different orientations of the magnetic field. The solid line represents a linear function αT with $\alpha = 283 \frac{\omega_c \tau}{\epsilon_F B} = 0.38$ mV K⁻² T⁻¹ (see text and Table I). Both this function and the low-temperature data are displayed in the inset as a ν/T vs T plot.



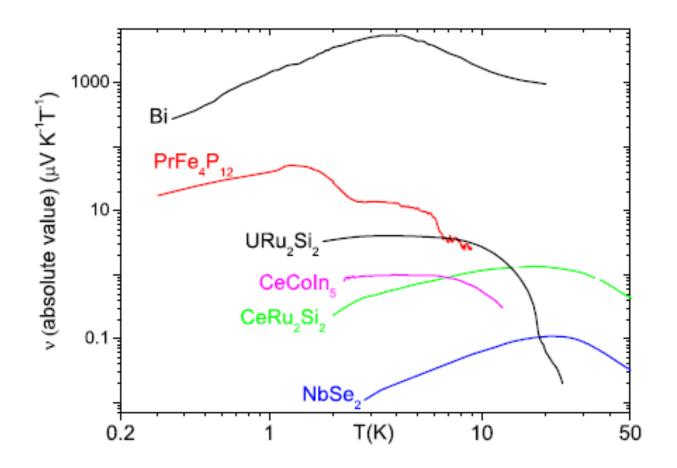


FIG. 3 (color online). The magnitude of the Nernst coefficient in bismuth compared to what is found in some other metals [4,5,7,8,12].

K. Behnia group



Giant Nernst Effect and Lock-In Currents at Magic Angles in (TMTSF)₂PF₆

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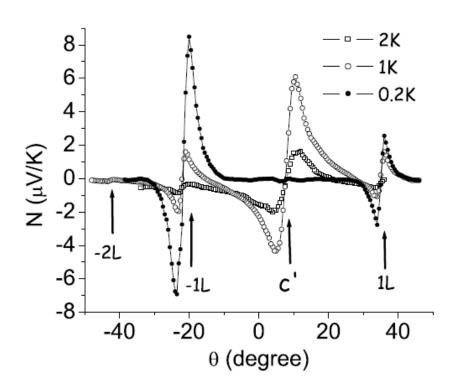


FIG. 2. Angular dependence of the Nernst signal at 2, 1, and 0.2 K, B = 7.5 T showing the rapid growth with lower temperature. We have used $N = \frac{S(B) - S(-B)}{2}$ as our definition for the Nernst signal.

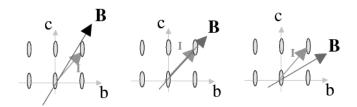


FIG. 4. Looking down the TMTSF chains. For the field near the $\mathbf{b} + \mathbf{c}$ direction the current flows only between the chains separated by $\mathbf{b} + \mathbf{c}$. The Lorentz force then produces a force along \mathbf{a} in the first figure, along $-\mathbf{a}$ in the last figure, and no force in the middle figure when the field and the current are parallel. Note that here we use an orthorhombic approximation.



Reading:

General

J. M. Ziman Principles of the theory of solids

Nernst effect in superconductors, review R. P. Huebener, Supercond. Sci. Technology 8, 189 (1995).

Nernst effect in exotic materials

K. Behnia, virtual lecture at https://www.nc/workshop/work_lectures.html

Review articles on physics of thermal and thermoelectric phenomena N. Hussey Adv. Phys. 51, 1685 (2002).

K. Behnia, D. Jaccard, J. Flouquet, JPCM 16, 5187 (2003)